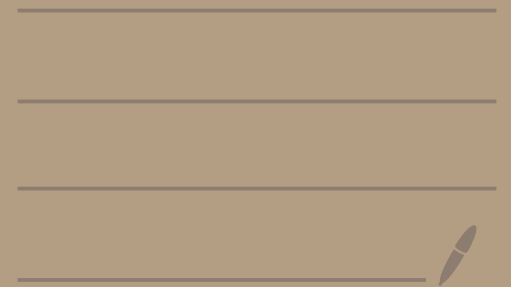


Lecture notes BIP Timisoara 2026



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INTRODUCTION

Relativistic hydrodynamics from Statistical Quantum Mechanics (SQM)

$$\hat{\rho} = \text{density operator} \quad \hat{\rho}^\dagger = \hat{\rho} \quad \text{Tr} \hat{\rho} = 1 \quad \hat{\rho}^2 \leq \hat{\rho}$$

$$T^{\mu\nu} \rightarrow \text{Tr}(\hat{\rho} \hat{T}^{\mu\nu}(x)) \quad \partial_\mu T^{\mu\nu} \rightarrow \text{Tr}(\hat{\rho} \partial_\mu \hat{T}^{\mu\nu}) = 0$$

How to obtain $\hat{\rho}$?

GLOBAL EQUILIBRIUM

Entropy in SQM

$$S = -\text{Tr}(\hat{\rho} \log \hat{\rho})$$

$\hat{\rho}$ can only depend on globally constant

max entropy (basic principle of S(Q)M)

10 constants P_0^μ $J_0^{\mu\nu}$ mean values + some charges Q_i

Each constant \leftrightarrow Lagrange multiplier

$$F[\hat{\rho}] = -\text{Tr}(\hat{\rho} \log \hat{\rho}) - b_{\mu} (\text{Tr}(\hat{\rho} \hat{P}^{\mu}) - P_0^{\mu}) + \underline{\omega}_{\mu\nu} (\text{Tr}(\hat{\rho} \hat{J}^{\mu\nu}) - J_0^{\mu\nu}) \\ + \sum_i \zeta_i (\text{Tr}(\hat{\rho} \hat{Q}_i) - Q_{0i}) + \lambda (\text{Tr}(\hat{\rho}) - 1)$$

$$\frac{\delta F}{\delta \hat{\rho}} = 0 \quad \text{Solution} \quad \hat{\rho} = \frac{e^{-b \cdot \hat{P} + \frac{1}{2} \underline{\omega} : \hat{J} + \sum_i \zeta_i \hat{Q}_i}}{Z}$$

$$Z = \text{Tr} (e^{-b \cdot \hat{P} + \frac{1}{2} \underline{\omega} : \hat{J} + \sum_i \zeta_i \hat{Q}_i})$$

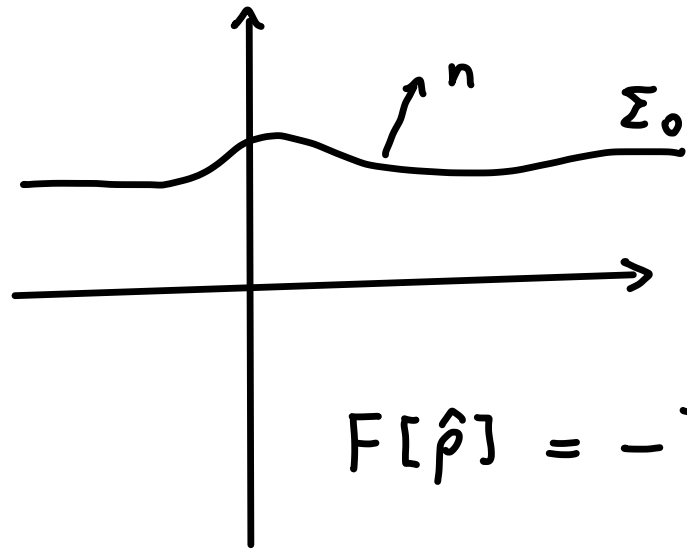
b = constant future-timelike vector

$\underline{\omega}$ = constant anti-symmetric tensor

Problem : prove it (T)

LOCAL EQUILIBRIUM

Maximize entropy with DENSITIES of conserved currents



$$n_\mu \text{Tr}(\hat{\rho} \hat{T}^{\mu\nu}(x)) = n_\mu T_{\text{given}}^{\mu\nu}(x) \text{ on } \Sigma_0$$

$$n_\mu \text{Tr}(\hat{\rho} \hat{j}^\mu(x)) = n_\mu j_{\text{given}}^\mu(x) \text{ on } \Sigma_0$$

$$F[\hat{\rho}] = -\text{Tr}(\hat{\rho} \log \hat{\rho}) - \int_{\Sigma_0} d\Sigma_\mu [\beta_\nu(x) (\text{Tr}(\hat{\rho} \hat{T}^{\mu\nu}(x)) - T_{\text{given}}^{\mu\nu})]$$

$$\zeta(x) [\text{Tr}(\hat{\rho} \hat{j}^\mu(x)) - j_{\text{given}}^\mu] + \lambda (\text{Tr}(\hat{\rho}) - 1)$$

β, ζ solutions
of the constraints

Solution:

$$\hat{\rho}_{\text{LE}}(\Sigma_0) = \exp \left[- \int_{\Sigma_0} d\Sigma_\mu \hat{T}^{\mu\nu} \beta_\nu - \zeta \hat{j}^\mu \right] / Z$$

Problem: prove that β_μ is a Killing vector at G.E. in G.R. (T)

Physical meaning of β and ζ ?

$$S = -\text{Tr}(\hat{\rho}_{LE} \log \hat{\rho}_{LE}) = \log Z + \int d\Sigma_\mu \beta_\nu \langle \hat{T}^{\mu\nu} \rangle_{LE} - \int d\Sigma_\mu \zeta \langle \hat{j}^\mu \rangle_{LE}$$

Suppose β is constant and so ζ

$$\begin{aligned} \Rightarrow S &= \log Z + \beta_\nu \int d\Sigma_\mu \langle \hat{T}^{\mu\nu} \rangle_{LE} - \zeta \int d\Sigma_\mu \langle \hat{j}^\mu \rangle_{LE} \\ &= \log Z + \beta_\nu \langle \hat{P}^\nu \rangle_{LE} - \zeta \langle \hat{Q} \rangle_{LE} \end{aligned}$$

Comparing with classical expression $TS = \log Z + u \cdot P - \mu Q$

$$\beta^\mu = \frac{1}{T} u^\mu \quad \zeta = \frac{\mu}{T} \quad \forall x!$$

Problem: prove that $\beta_\mu = b_\mu + \omega_{\mu\nu} x^\nu$ at global equilibrium (T)

Problem: different cases for b, ω ; Tolman's law (T)

Calculate $T_{LE}^{\mu\nu}(x) = \text{Tr}(\hat{\rho}_{LE} \hat{T}^{\mu\nu}(x))$

$$\text{Tr}(\hat{\rho}_{LE} \hat{T}^{\mu\nu}(x)) = \frac{1}{Z_{LE}} \text{Tr} \left(e^{-\int_{\Sigma_0} d\Sigma_\rho (\hat{T}^{\rho\sigma} \beta_\sigma - \zeta \hat{J}^\rho)} \hat{T}^{\mu\nu}(x) \right)$$

Idea: mean value in x mostly determined by β, ζ around x

$$\begin{aligned} e^{-\int_{\Sigma_0} d\Sigma_\rho(y) \hat{T}^{\rho\sigma}(y) \beta_\sigma(y)} &\cong e^{-\int_{\Sigma_0} d\Sigma_\rho(y) \hat{T}^{\rho\sigma}(y) [\beta_\sigma(x) + \partial_\lambda \beta_\sigma(x) (y-x)^\lambda + \dots]} \\ &= e^{-\beta_\sigma(x) \underbrace{\int_{\Sigma_0} d\Sigma_\rho(y) \hat{T}^{\rho\sigma}(y)}_{\hat{P}^\sigma} + \partial_\lambda \beta_\sigma(x) \int_{\Sigma_0} d\Sigma_\rho(y) (y-x)^\lambda \hat{T}^{\rho\sigma} + \dots} \\ &= e^{-\beta_\sigma(x) \hat{P}^\sigma + \text{corrections}} \end{aligned}$$

At the lowest order

$$T_{LE}^{\mu\nu}(x) \cong \text{Tr} \left(\frac{e^{-\beta(x) \cdot \hat{P}} \hat{T}^{\mu\nu}(x)}{Z_0} \right)$$

$$\rightarrow T_{LE}^{\mu\nu}(x) \cong A \beta^\mu \beta^\nu + B g^{\mu\nu}$$

equivalent to $(\rho + p) u^\mu u^\nu - p g^{\mu\nu}$

Proof \rightarrow tutorial 1

Constitutive equations at LE

$$T_{LE}^{\mu\nu}(x) = T_{ideal}^{\mu\nu} + \text{quantum corrections}$$

ENTROPY CURRENT

$$1) S = -\text{Tr}(\hat{\rho} \log \hat{\rho}) \xrightarrow{\frac{d\hat{\rho}}{dt}} -\text{Tr}(\hat{\rho}_{LE} \log \hat{\rho}_{LE}) \quad \frac{d\hat{\rho}_{LE}}{dt} \neq 0$$

maximized
↓ constrained ignorance

2) Extensivity

$$\text{Let } \hat{\rho}(\lambda) = \frac{1}{Z(\lambda)} \exp\left[-\lambda \int d\Sigma_\mu \hat{T}^{\mu\nu} \beta_\nu\right] \quad (\zeta=0 \text{ for simplicity}) \quad \hat{\rho} \equiv \hat{\rho}(1)$$

$$Z(\lambda) = \text{Tr}\left(\exp\left[-\lambda \int d\Sigma_\mu \hat{T}^{\mu\nu} \beta_\nu\right]\right) \Rightarrow \frac{\partial \log Z}{\partial \lambda} = \frac{1}{Z(\lambda)} \frac{\partial}{\partial \lambda} \text{Tr}\left(\exp\left[-\lambda \int d\Sigma_\mu \hat{T}^{\mu\nu} \beta_\nu\right]\right)$$

$$= \frac{\text{Tr}\left(\exp\left[-\lambda \int d\Sigma_\mu \hat{T}^{\mu\nu} \beta_\nu\right] (-1) \int d\Sigma_\mu \hat{T}^{\mu\nu} \beta_\nu\right)}{Z(\lambda)} = - \int d\Sigma_\mu \langle \hat{T}^{\mu\nu} \rangle(\lambda) \beta_\nu$$

$\langle \hat{T}^{\mu\nu} \rangle(\lambda)$
mean value with λ

Integrating between 1 and λ_0

$$\int_1^{\lambda_0} d\lambda \frac{\partial}{\partial \lambda} \log Z(\lambda) = \log Z(\lambda_0) - \log Z = - \int_1^{\lambda_0} d\lambda \int d\Sigma_\mu \langle \hat{T}^{\mu\nu} \rangle(\lambda) \beta_\nu = \int d\Sigma_\mu \left(\int_{\lambda_0}^1 d\lambda \langle \hat{T}^{\mu\nu} \rangle(\lambda) \beta_\nu \right)$$

If λ_0 such that $\log Z(\lambda_0) = 0$

$$\Rightarrow \log Z = \int d\Sigma_\mu \phi^\mu \quad \text{with} \quad \phi^\mu = \int_1^{\lambda_0} d\lambda \langle \hat{T}^{\mu\nu} \rangle(\lambda) \beta_\nu$$

Does it exist?

$$\text{Suppose } \lambda_0 = +\infty \Rightarrow \lim_{\lambda_0 \rightarrow \infty} \frac{e^{-\lambda_0 \int d\Sigma_\mu \hat{T}^{\mu\nu} \beta_\nu}}{Z(\lambda)} = |0\rangle\langle 0|$$

$$|0\rangle = \text{LLS of } \int d\Sigma_\mu \hat{T}^{\mu\nu} \beta_\nu$$

Project: calculate ϕ^μ for $\hat{\rho} = e^{-b \cdot \hat{P}} / Z$

Entropy current

$$S = \int d\Sigma_\mu \phi^\mu + T^{\mu\nu}_{LE} \beta_\nu - \int j^\mu_{LE}$$

$$S = \int d\Sigma_\mu \phi^\mu + T^{\mu\nu} \beta_\nu - \int j^\mu$$

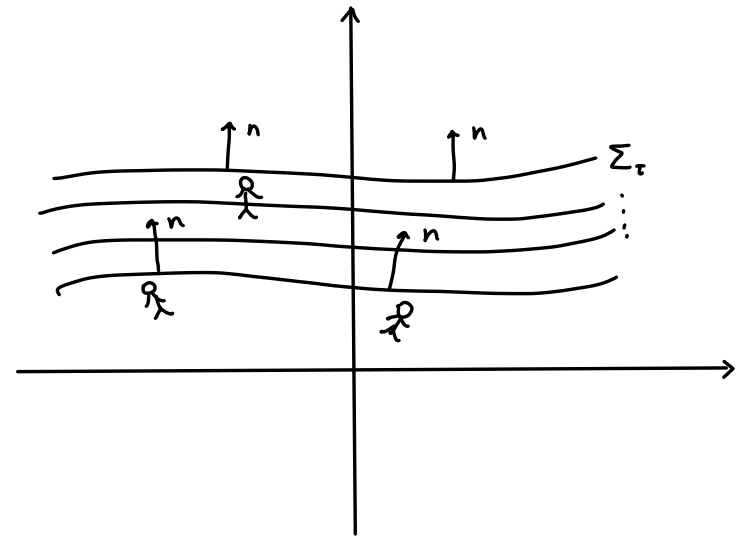
$$\Rightarrow S^\mu = \phi^\mu + T^{\mu\nu} \beta_\nu - j^\mu$$

is it independent of Σ (better: the foliation?)

No, because $\int (n)$ and $\beta(n)$ both depend on the foliation.

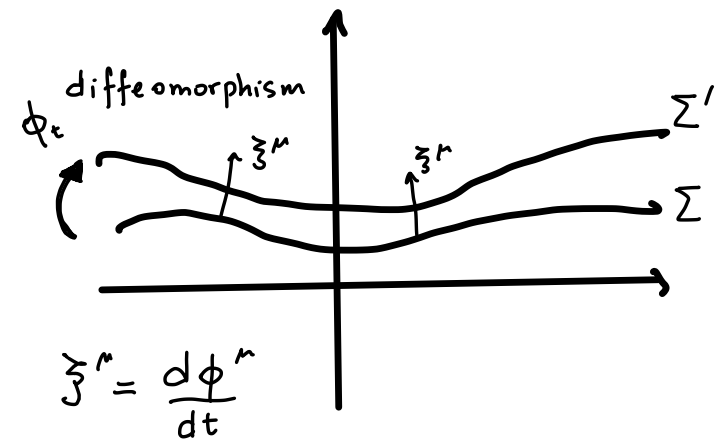
It can be made independent only if there was a privileged foliation or at global equilibrium.

$$n_\mu T^{\mu\nu} = n_\mu T^{\mu\nu}_{LE}$$
$$n_\mu j^\mu = n_\mu j^\mu_{LE}$$



Entropy production

$$\frac{\delta}{\delta \Sigma} \left(\int_{\Sigma} d\Sigma_r V^r \right)_{\xi} = \frac{1}{2} \int d\tilde{S}_{\mu\nu} \xi^\nu V^\mu + \int d\Sigma_r \xi^r \nabla \cdot V$$



$$Z_{\phi_t(\Sigma)} - Z_{\Sigma} = \int_{\phi_t(\Sigma)} d\Sigma_r \phi^r - \int d\Sigma_r \phi^r \implies \mathcal{L}_{\xi} \int d\Sigma_r \phi^r \text{ domain derivative} = \int d\Sigma_r \xi^r \nabla \cdot \phi \text{ if } \int d\tilde{S} = 0$$

On the other hand:

$$Z_{\phi_t(\Sigma)} = \text{Tr} \left(e^{-\int_{\phi_t(\Sigma)} d\Sigma_r \hat{T}^{\mu\nu} \beta_\nu} \right) \cong \text{Tr} \left(e^{-\int d\Sigma_r \hat{T}^{\mu\nu} \beta_\nu - \varepsilon \int d\Sigma_r \xi^r \nabla \cdot (\hat{T}^{\mu\nu} \beta_\nu)} \right) \approx \int d\tilde{S} = 0$$

$$\cong Z - \varepsilon \text{Tr} \left(e^{-\int \dots} \int d\Sigma \cdot \xi \nabla_\mu (\hat{T}^{\mu\nu} \beta_\nu) \right)$$

whence

$$\lim_{\varepsilon \rightarrow 0} \frac{Z_\varepsilon - Z}{\varepsilon} = \mathcal{L}_{\xi} (\log Z) = - \int d\Sigma \cdot \xi \langle \hat{T}^{\mu\nu} \rangle_{\varepsilon} \nabla_\mu \beta_\nu$$

$$\Rightarrow \int d\Sigma_\mu \phi^\mu = \int d\Sigma \cdot \xi \nabla \cdot \phi = - \int d\Sigma \cdot \xi \langle \hat{T}^{\mu\nu} \rangle_{LE} \nabla_\mu \beta_\nu$$

being ξ arbitrary $\Rightarrow \nabla \cdot \phi = - \langle \hat{T}^{\mu\nu} \rangle_{LE} \nabla_\mu \beta_\nu$

and, with current, $\nabla \cdot \phi = - \langle \hat{T}^{\mu\nu} \rangle_{LE} \nabla_\mu \beta_\nu + \langle \hat{j}^\mu \rangle_{LE} \nabla_\mu \zeta$

$$\Rightarrow \nabla \cdot s = \nabla \cdot (\phi + T^{\mu\nu} \beta_\nu - j^\mu) = \nabla \cdot \phi + \cancel{\nabla_\mu T^{\mu\nu}} \beta_\nu + T^{\mu\nu} \nabla_\mu \beta_\nu - \cancel{\nabla_\mu j^\mu} \zeta$$

$$- \nabla \zeta \cdot j = - T^{\mu\nu}_{LE} \nabla_\mu \beta_\nu + j^\mu_{LE} \nabla_\mu \zeta + T^{\mu\nu} \nabla_\mu \beta_\nu - \nabla_\mu \zeta j^\mu \Rightarrow$$

$$\nabla \cdot s = (T^{\mu\nu} - T^{\mu\nu}_{LE}) \nabla_\mu \beta_\nu - (j^\mu - j^\mu_{LE}) \nabla_\mu \zeta$$